

# F-K Domain Wavefield Continuation with Arbitrary Velocities

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SEP120 page: 311-317

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# Motivation

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- Mixed-domain methods interpolate wavefields to account for laterally varying velocities. Can we interpolate phases instead?

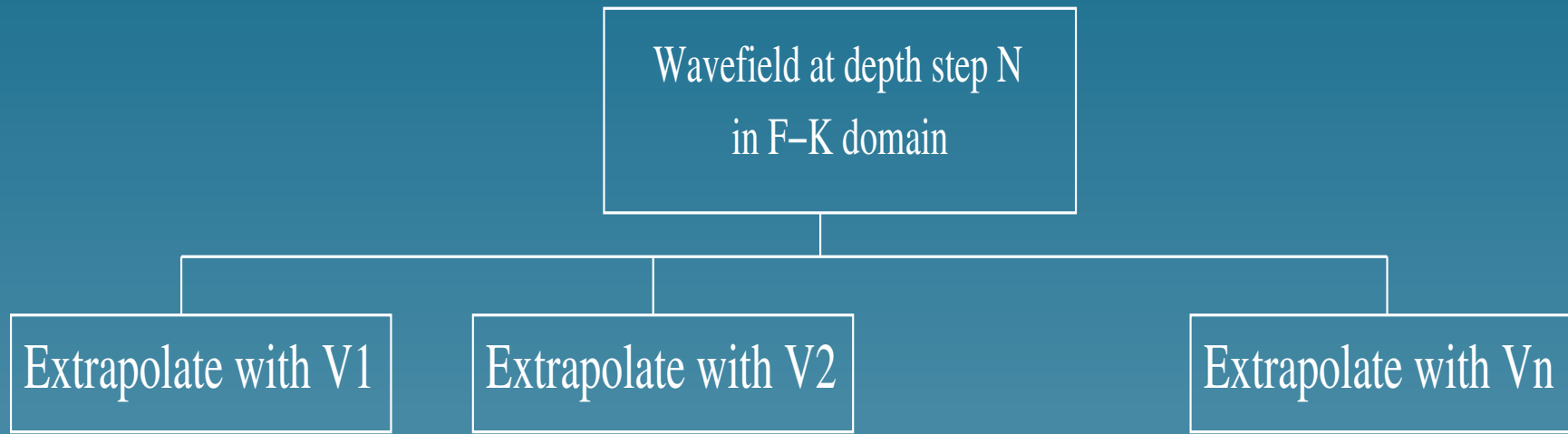
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- Mixed-domain methods interpolate wavefields to account for laterally varying velocities. Can we interpolate phases instead?
- Interpolating phases will allow arbitrary velocity variations and a faster and simpler algorithm

# Overview of PSPI (I)

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# Overview of PSPI (II)

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In  $\omega$ - $\mathbf{k}$  space:

$$\mathbf{W}_l^{N+1} = \mathbf{W}^N e^{ik_{z_l}\Delta z}$$

where:

$\mathbf{W}^N$ : Wavefield at depth  $N$

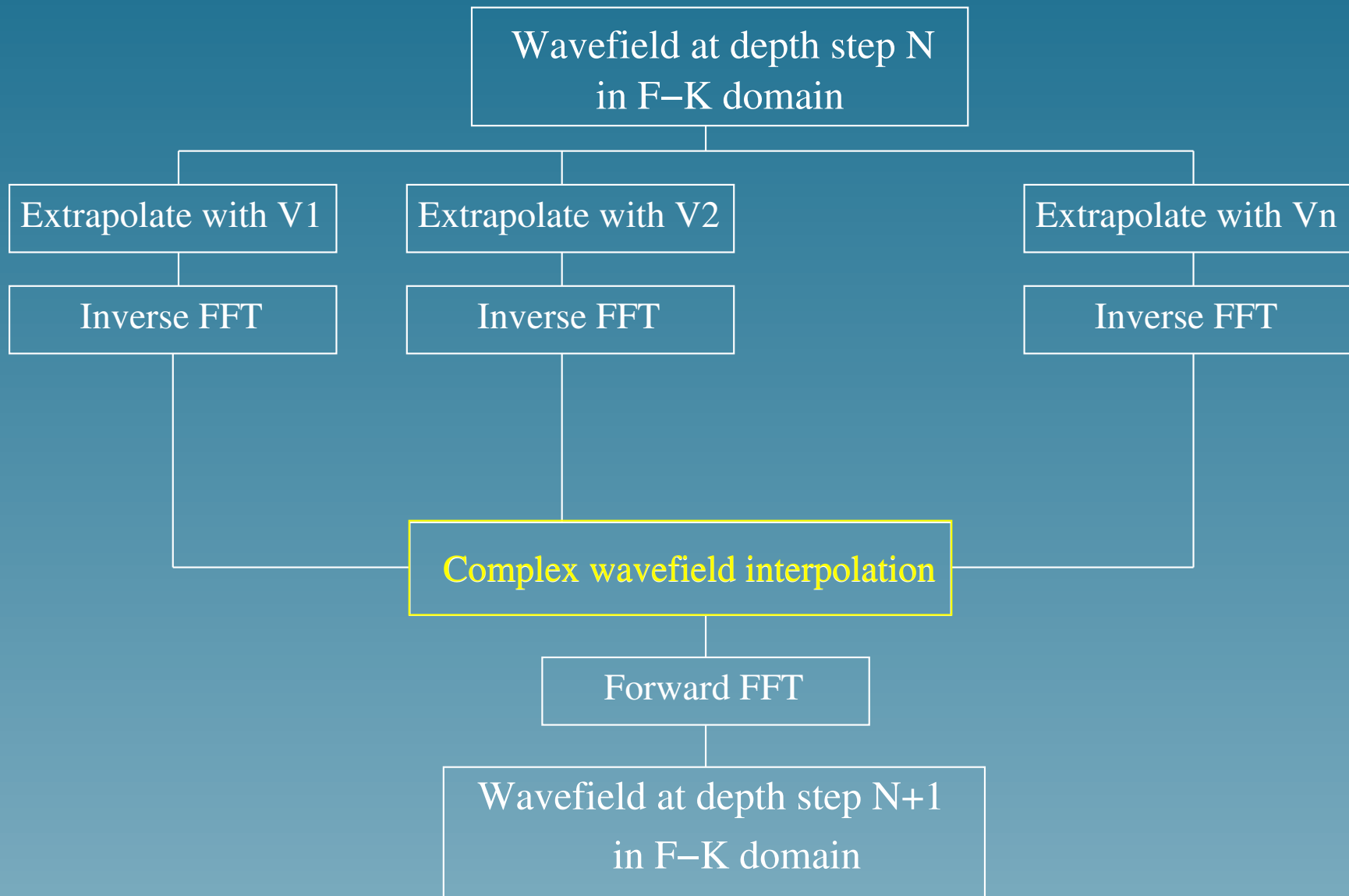
$V_l$ :  $l$ -th Reference velocity

$\mathbf{W}_l^{N+1}$ : Wavefield at depth  $N + 1$  continued with  $V_l$

and

$$k_{z_l} = \sqrt{\frac{\omega^2}{V_l^2} - |\mathbf{k}|^2} \quad \text{is the dispersion relation}$$

# Overview of PSPI (III)



# Overview of PSPI (IV)

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In  $\omega$ - $\mathbf{x}$  space:

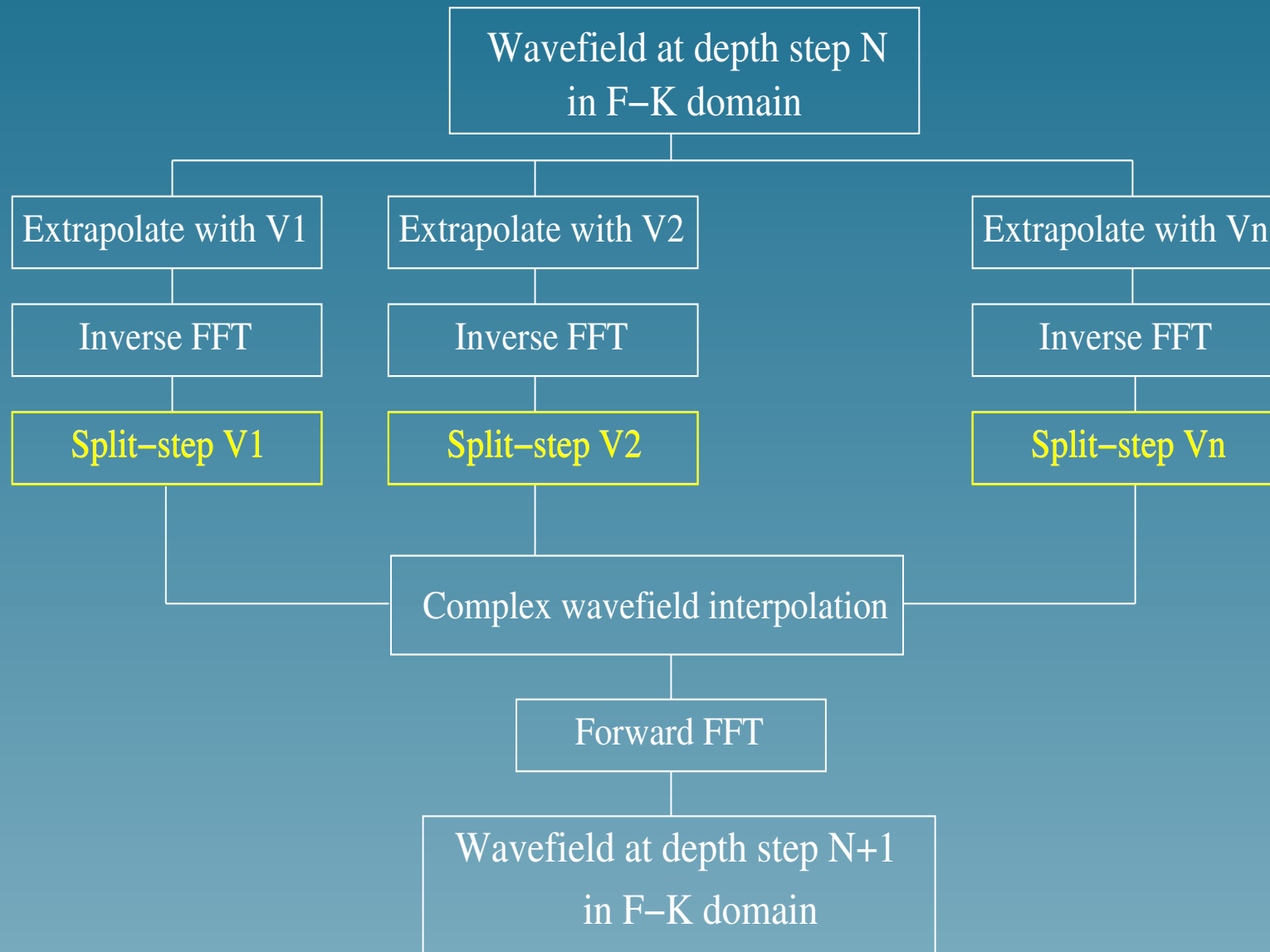
$$\mathbf{w}^{N+1}(j) = \sum_{l=1}^{nv} \sigma_l \mathbf{w}_l^N(j)$$

where:

$\sigma_l$ : interpolation factor

$\mathbf{w}_l^{N+1}(j)$ : wavefield in  $\omega$ - $\mathbf{x}$  at location  $j$

# Overview of Extended Split-step





# Overview of Split-step Correction

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The split-step correction is given by:

$$e^{i\left(\frac{\omega}{V} - \frac{\omega}{V_l}\right)\Delta z},$$

where  $V$  is the true velocity is applied before the interpolation and is intended to compensate, to a first order, for the difference between  $V$  and  $V_l$ .

# The Idea of Interpolating Phases (I)

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In Phase shift extrapolation for  $V(z)$ :

$$\mathbf{W}^{N+1} = \mathbf{W}^N e^{i\theta z}$$

In  $V(\mathbf{x}, z)$  find an “equivalent” phase such that:

$$\mathbf{W}^{N+1} = \mathbf{W}^N e^{i\theta_{z\text{eq}}}$$

# The Idea of Interpolating phases (II)

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For any two complex numbers  $z_1 = A\theta_1$  and  $z_2 = A\theta_2$ :

$$\Phi\left(\frac{z_1 + z_2}{2}\right) = \frac{\theta_1 + \theta_2}{2}$$

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and

$$\text{Amp}\left(\frac{z_1 + z_2}{2}\right) = \frac{A}{\sqrt{2}}\sqrt{1 + \cos(\theta_2 - \theta_1)} \neq A$$

# The Proposed Algorithm

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Wavefield Extrapolation with Arbitrary Velocities in  $\omega$ - $\mathbf{K}$

# Basic Concept

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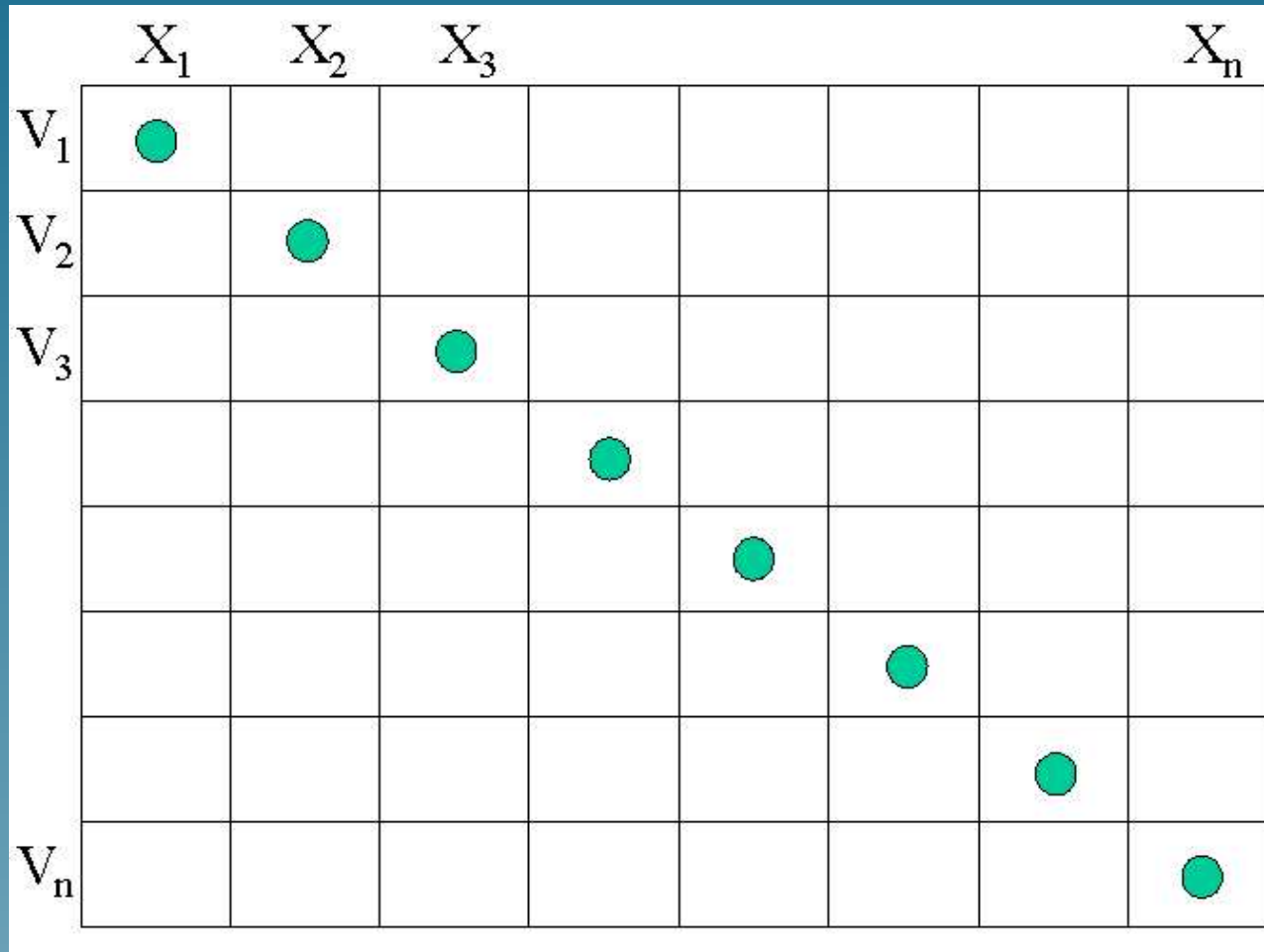
- Assume that as many reference velocities as spatial locations are used at each depth step.
- $nv = n\mathbf{x}$ .
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# Basic Concept

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- Assume that as many reference velocities as spatial locations are used at each depth step.
- $n_v = n_x$ .
- No split-step correction is required.
- No need for high-order approximation of the dispersion relation.
- Wavefield interpolation is replaced by selection.

# Wavefield Selection



Each row is a wavefield extrapolated with the indicated velocity.

# Extrapolated Wavefield in $\omega$ - $\mathbf{x}$

---

The selection process in the  $\omega$ - $\mathbf{x}$  is given by:

$$\mathbf{w}^{N+1}(j) = \sum_{l=1}^{nv} \mathbf{w}_l^{N+1}(j) \delta_{lj}$$

where

$\mathbf{w}_l^{N+1}$ :  $l$ th row in the array of extrapolated wavefields.

$\delta_{lj}$ : Kronecker delta to select the  $j = l$  component.

# Extrapolated Wavefield in $\omega$ - $\mathbf{K}$

---

The equivalent equation in the  $\omega$ - $\mathbf{K}$  domain is:

$$\mathbf{W}^{N+1} = \sum_{l=1}^{nv} \mathbf{W}_l^{N+1} \otimes e^{-ik_x \Delta x_l}$$

where

$$\Delta x_l = (l - 1) \Delta x / nx$$

$\otimes$ : circular convolution

One spatial index is used to simplify the notation.

# Extrapolated Wavefield in $\omega$ - $\mathbf{K}$ (II)

---

The extrapolated wavefield in  $\omega$ - $\mathbf{K}$  is then:

$$\mathbf{W}^{N+1}(j) = \sum_{l=1}^{nv} \sum_{m=\langle nx \rangle} \mathbf{W}_l^{N+1}(m) e^{-ik_x(j-m)\Delta x_l}$$

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Replacing  $\mathbf{W}_l^{N+1}$  in terms of  $\mathbf{W}^N$  and rearranging terms:

$$\mathbf{W}^{N+1}(j) = \sum_{m=1}^{nx} \mathbf{W}^N(m) \sum_{l=1}^{nv} e^{-ik_{z_l}(m)\Delta z + k_x(\tilde{m}_j)\Delta x_l}$$

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where  $\tilde{m}_j = \text{mod}(j - m, nx)$



# Extrapolated Wavefield in $\omega$ - $\mathbf{K}$ (III)

---

Written as a dot product:

$$\mathbf{W}^{N+1}(j) = \sum_{m=1}^{nx} \mathbf{W}^N(m) \mathbf{f}_j(m) = \mathbf{W}^N \cdot \mathbf{f}_j.$$

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The vector  $\mathbf{f}_j$  is independent of the data and contains the velocity information:

$$\mathbf{f}_j = \sum_{l=1}^{nv} e^{-ik_{z_l}(m)\Delta z + k_x(\tilde{m}_j)\Delta x_l}.$$

# Practical Implementation (I)

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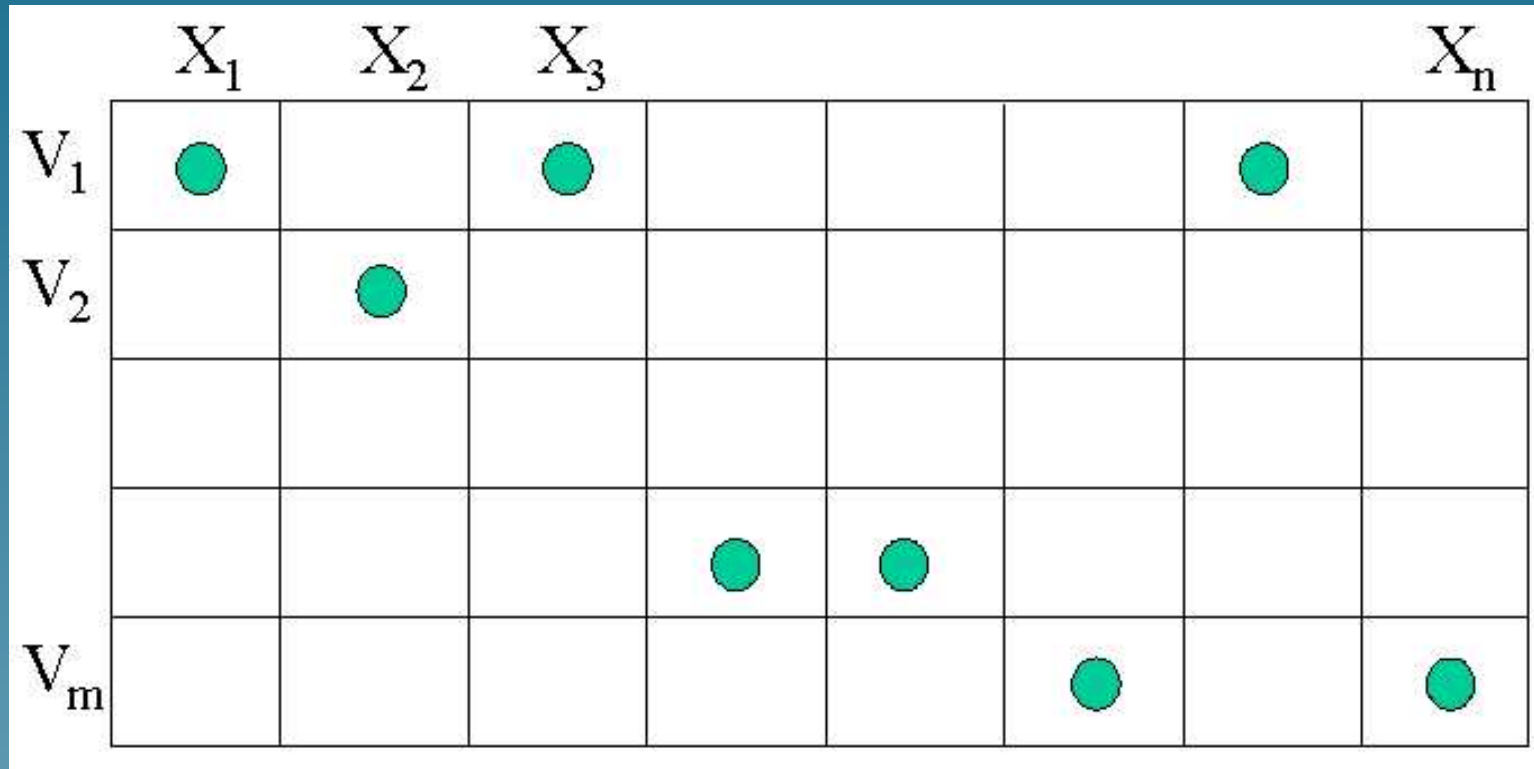
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# Practical Implementation (II)

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- The algorithm can be made essentially quadratic by realizing that:
  - ★ Velocities can be binned to within their assumed accuracy.
  - ★ The vertical wavenumber can be precomputed, since it does not depend on velocity.

# Modified Wavefield Selection



This time each row is an extrapolated wavefield with the indicated **binned velocity**.

# Modified Wavefield in $\omega$ - $\mathbf{x}$

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The selection process to calculate the wavefield in  $\omega$ - $\mathbf{x}$  is now:

$$\mathbf{w}^{N+1} = \sum_{l=1}^{nv} \mathbf{w}_l^{N+1} \sum_p \delta_{pl}.$$



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$l$ : velocity index.

$p$ : index to select spatial locations with the same velocity.

# Modified Wavefield in $\omega$ - $\mathbf{K}$

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In the  $\omega$ - $\mathbf{K}$  domain:

$$\mathbf{W}^{N+1}(j) = \sum_{m=1}^{nx} \mathbf{W}^N(m) \sum_{l=1}^{nv} \left( e^{-ik_{z_l}(m)\Delta z} \sum_p e^{-ik_x(\tilde{m}_j)\Delta x_p} \right)$$

# Extrapolated Wavefield

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Conceptually, the result is the same that we obtained before:

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only this time the vector  $\mathbf{f}_j$  is given by

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# Remarks

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- No significant approximations have been made.
- The cost comes from considering every wavefield trace in the computation of every other one.

# Speculative Ideas on Improving Efficiency

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- Subsample the wavefield used for the computation of each wavefield trace at the next depth step.



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- Compute only a subsampled version of the wavefield and interpolate.
- Interpolate phases two by two.

# Subsample the Input Wavefield

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- Subsampling in wavenumber domain implies windowing in the space domain.
- May be a better approximation at shallow than at deeper depths.

# Subsample the Computed Wavefield

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- Linearly interpolate for the wavefield traces not computed.
- This implies that the wavefield is somewhat smooth in the spatial direction.

# Interpolating Phases Two-by-Two

---

For any two phases  $\theta_1$  and  $\theta_2$ :

$$\Phi \left( \frac{e^{i\theta_1} + e^{i\theta_2}}{2} \right) = \frac{\theta_1 + \theta_2}{2}$$

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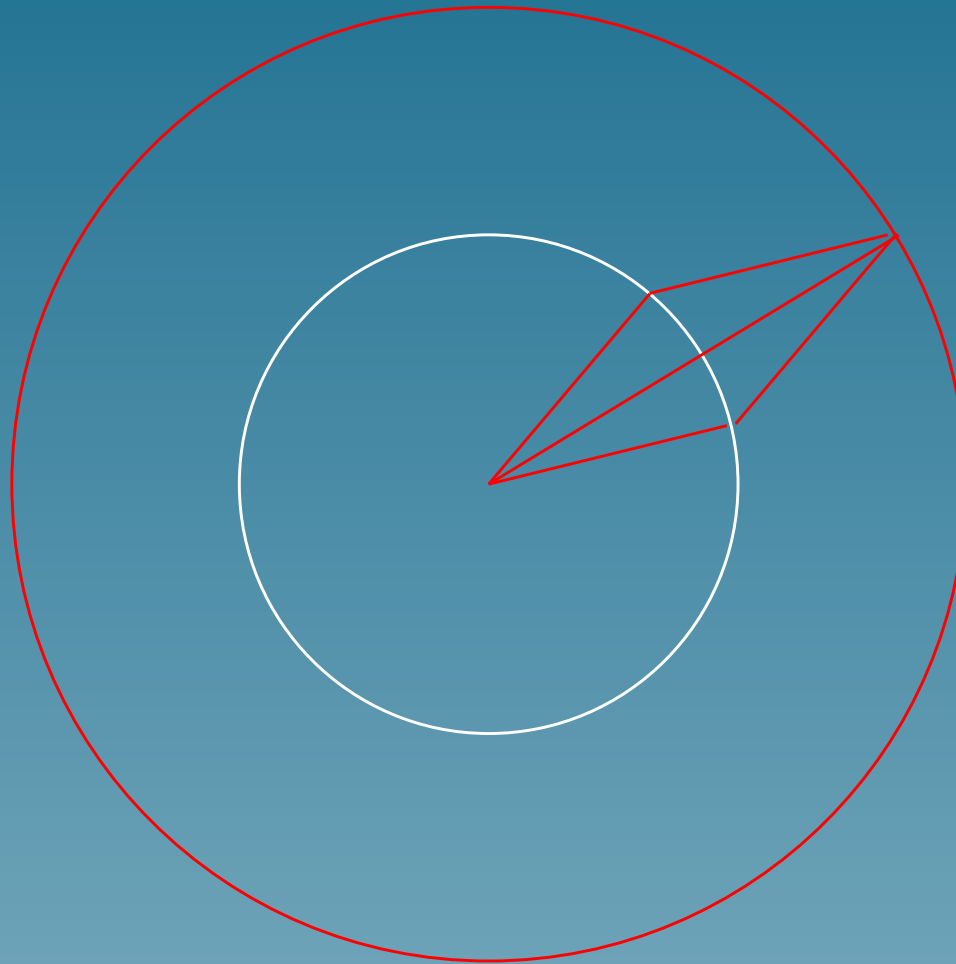
But

$$\text{Amp} \left( \frac{e^{i\theta_1} + e^{i\theta_2}}{2} \right) = \frac{1}{\sqrt{2}} \sqrt{1 + \cos(\theta_2 - \theta_1)} \neq 1$$

The question is: can we pair-up the sum of exponentials such that the amplitude term becomes a normalization?

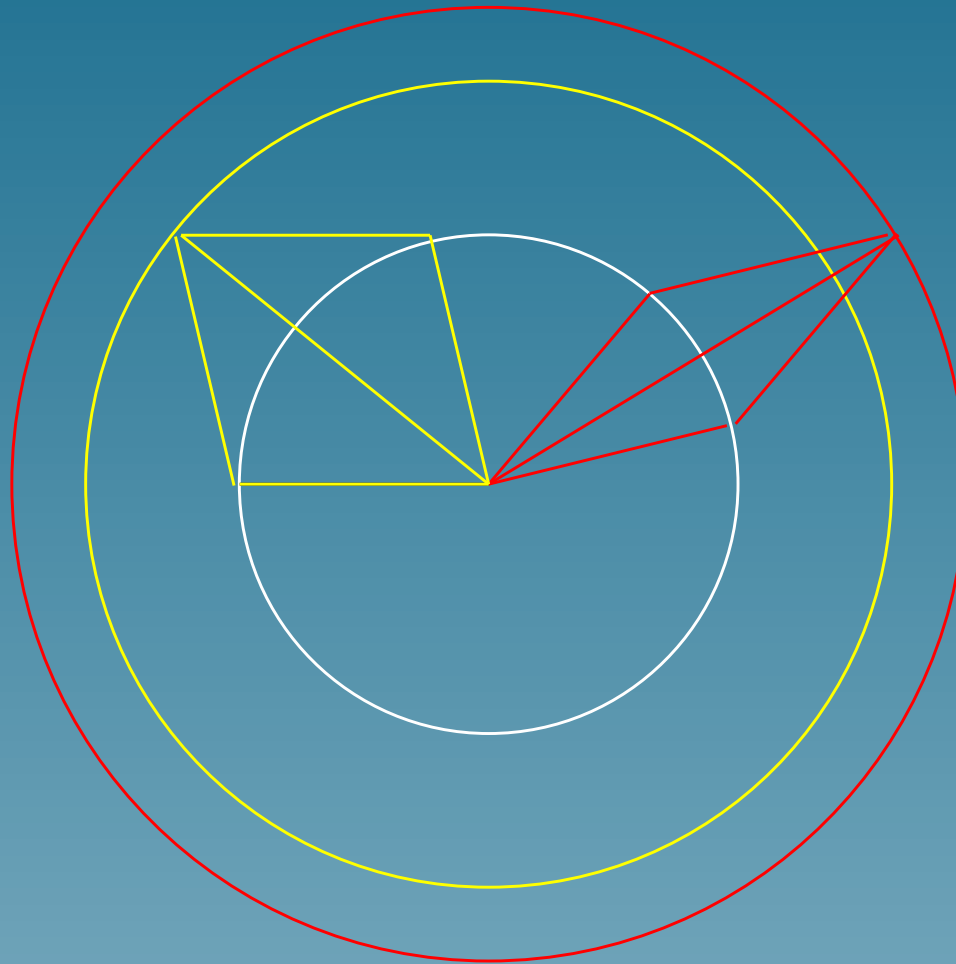
# Interpolating Phases Sketch

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- It is possible to do F-K wavefield continuation with arbitrary spatial velocity variations.
- The resulting algorithm is quadratic in the spatial dimensions so needs to be made more efficient.
- We have given some untested ideas on how to overcome the high cost. Testing those ideas is the next step.