

FREQUENCY DOMAIN IMPLEMENTATION OF
SLANT-MIDPOINT IMAGING

Richard Ottolini

Slant-midpoint stacks are as easily migrated in the frequency domain as CDP stacks. Slant-midpoint imaging enhances selected dip ranges (other papers in this report) and assists lateral velocity estimation [0]. This paper gives the frequency domain migration and diffraction transfer functions for CDP stacked, slant stacked, unstacked data images in common midpoint coordinates. The peculiar shapes of the computed slant-midpoint point scatterer response are theoretically verified by Clayton [1]. A fast computer program for modeling or migrating slant-midpoint images is included.

A wave equation transfer function may be translated entirely into the frequency domain by Stolt's method [2,3]. Consider the one way wave equation solution in *common-midpoint-stacked* coordinates

$$P_{\text{mig}}(k_x, z, \omega) = P_{\text{dif}}(k_x, z_0, \omega) \exp iz \frac{\omega}{v} \sqrt{1 - \frac{v^2 k_x^2}{\omega^2}}$$

where k_x is midpoint wave number

$$k_z = \frac{\omega}{v} \sqrt{1 - \frac{v^2 k_x^2}{\omega^2}} \quad \text{is depth wave number}$$

The inverse fourier transform is

$$P_{\text{mig}}(x, z, t) = \iint_{-\infty}^{\infty} d\omega dk_x e^{ik_x x - i\omega t} P_{\text{dif}}(k_x, z_0, \omega) e^{i \frac{\omega}{v} \sqrt{1 - \frac{v^2 k_x^2}{\omega^2}} z}$$

where k_x is midpoint wave number as before

$$H = \frac{\omega}{2v} k_h + pv \text{ is slant stack parameter (constant)}$$

k_h is half offset wave number

$$k_z = \frac{\omega}{v} \left[\sqrt{\dots} + \sqrt{\dots} \right] \text{ is depth wave number}$$

Migration is

$$P_{\text{mig}}(k_x, k_z) = \frac{v}{2} \left[\frac{k_z^2(1-H^2) - 2k_x^2 k_z^2 H^2 - k_x^4 H^4}{(k_x^2 + k_z^2)^{1/2} (k_z^2(1-H^2) - k_x^2 H^2)^{3/2}} \right]$$

$$\bullet P_{\text{diff}} \left(k_x, \omega = \frac{vk_z}{2} \sqrt{\frac{k_x^2 + k_z^2}{k_z^2(1-H^2) - k_x^2 H^2}} \right)$$

Diffraction is

$$P_{\text{dif}}(k_x, \omega) = \left[\frac{\frac{1}{v} (1-H^2) + \frac{1}{4} k_x H}{\sqrt{\omega^2(1-H^2) + \frac{\omega}{2v} k_x H - \frac{1}{4} v^2 k_x^2}} + \frac{\frac{1}{v} (1-H^2) - \frac{1}{4} k_x H}{\sqrt{\omega^2(1-H^2) - \frac{\omega}{2v} k_x H - \frac{1}{4} v^2 k_x^2}} \right]$$

$$\bullet P_{\text{mig}} \left(k_x, k_z = \left[\sqrt{\frac{\omega^2}{v^2} (1-H^2) + \frac{\omega}{2v} k_x H - \frac{1}{4} k_x^2} + \sqrt{\frac{\omega^2}{v^2} (1-H^2) - \frac{\omega}{2v} k_x H - \frac{1}{4} k_x^2} \right] \right)$$

Unstacked-midpoint transfer functions come from the wave equation solution given as Stolt's [3] equation 57 or Claerbout's [4] equation 7. Note they contain a third wave number dimension.

$$P_{\text{mig}}(k_x, k_h, z, \omega) = P_{\text{dif}}(k_x, k_h, z_0, \omega) \exp iz \frac{\omega}{v} \left(\sqrt{1 - \frac{(k_x - k_h)^2 v^2}{4\omega^2}} + \sqrt{1 - \frac{(k_x + k_h)^2 v^2}{4\omega^2}} \right)$$

Migration is

$$P_{\text{mig}}(k_x, k_h, k_z) = \frac{v}{z} \frac{k_x^2 k_h^2}{k_z^2} \sqrt{k_x^2 + k_h^2 + k_z^2 + \frac{k_x^2 k_h^2}{k_z^2}}$$

$$\bullet P_{\text{dif}}(k_x, k_h, \omega = \frac{v}{2} \sqrt{k_x^2 + k_h^2 + \frac{k_x^2 k_h^2}{k_z^2}})$$

Diffraction is

$$P_{\text{dif}}(k_x, k_h, \omega) = \frac{\omega}{v} \left[\frac{1}{\sqrt{\omega^2 - \frac{v}{4} (k_x^2 + k_h^2 - 2k_x k_h)}} - \frac{1}{\sqrt{\omega^2 - \frac{v}{4} (k_x^2 + k_h^2 + 2k_x k_h)}} \right]$$

$$\bullet P_{\text{mig}} \left[k_x, k_h, k_z = \left(\sqrt{\frac{\omega^2}{v^2} - \frac{1}{4} (k_x^2 + k_h^2 - 2k_x k_h)} + \sqrt{\frac{\omega^2}{v^2} - \frac{1}{4} (k_x^2 + k_h^2 + 2k_x k_h)} \right) \right]$$

Note the slant-midpoint-stacked and unstacked-midpoint transfer functions have the mathematical features

- (1) When H or k_h are zero they reduce to the common-midpoint-stacked case.
- (2) The mapping function portions (and obliquity functions) contain singularities. Nearby the singularities, the mapping functions try to select data from the nonexistent higher-than-nyquist-frequencies. Such locations may be safely zeroed out.
- (3) At low frequencies the mapping function may become complex valued before it becomes imaginary. Complex and imaginary

values are attenuated and evanescent energy respectively.

Though there is no computational objection to including attenuated energy, empirical results show severe distortion of the point scattering response. Attenuated energy should be zeroed out.

Documented Fortran programs for slant-midpoint migration and diffraction are attached. Minor modifications are necessary for the common-midpoint-stacked and unstacked-midpoint versions.

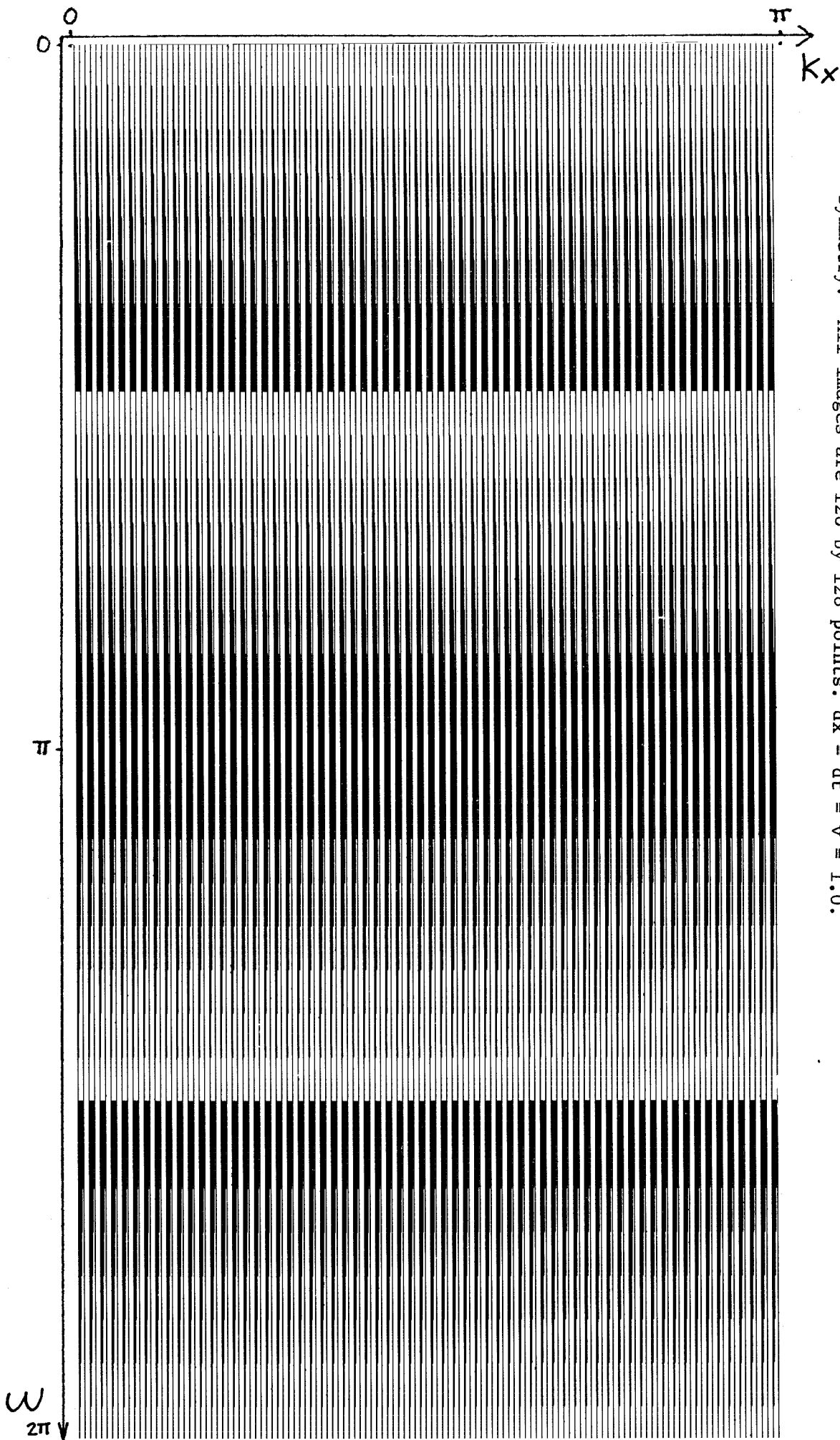
Figures 1 through 7 illustrate the behavior of the mapping functions on synthetic frequency domain images (k_x, ω and k_x, k_z spaces). As you recall, the mapping function is the portion of the transfer function which moves frequency values to different positions when performing frequency domain migration and diffraction. For these synthetic examples we will ignore the effects of the obliquity function, which is just a scale factor, by setting it to unity. Because positive k_x values are complex conjugates of negative k_x values, only half of the fourier images are displayed in the figures. Figure 1 is the k_x, ω source image for migrations and consists of straight lines of constant ω . Figure 2 is the slant midpoint migrated image for $H = 0$. Slant-midpoint for $H = 0$ and unstack at $kh = 0$ are equivalent to CDP migration. Since the CDP migration mapping function is an equation for a circle, the lines in figure 1 are transformed into the circles of figure 2. The zero regions of figure 2 represent attempts to map from non existant high ω frequencies. Figure 2 was used as the k_x, k_z source image for diffractions. Figure 3 is a slant-midpoint $H=0$ diffraction mapping where the circles of figure 2 have been mapped back into lines. The region labeled evanescent tried to map from nonexistent, imaginary k_z frequencies. Figures 4 and 5 show the slant-midpoint migration and diffraction mappings for ray parameter velocity product $H = pv = .3$. Non zero H values no longer map lines into circles and vice-versa. Additional regions of nonexistent and attenuated energy are introduced. Figure 6 is the same diffraction mapping for $H = .3$ as figure 5 adding the attenuated energy. Figure 7 is the unstacked-midpoint migration mapping for cross section $kh = .3$. Superficially, this mapping resembles the slant mid-point migration mapping of figure 4. However, the low k_x frequency content is considerably different.

Figures 8 and 9 illustrate the slant-midpoint point scatterer response. Figure 8 illustrates a suite of migrations and diffractions for different H values. Slant parameter $H=0$ of course generates the familiar common-midpoint migration semi-circle and diffraction hyperbola. Observe for non-zero H the migration responses no longer intersect the surfaces perpendicularly. This is due to the limited dip bandwidth of slant-midpoint stacking [1]. Figure 9 shows the disastrous effect of including attenuation energy during diffraction. Scattering responses for unstacked midpoint k_h cross sections are not very enlightening.

References

- [0] Lynn, W. S., "RMS Velocity Estimation in Laterally Varying Media," SEP 14
- [1] Clayton, R., "Migration in Midpoint Coordinates," SEP 14
- [2] Lynn, W. S., "Implementing f-k Migration and Diffraction," SEP 11
- [3] Stolt, R. H., "Migration by Fourier Transform," *Geophysics*, 43:23, Feb., 1978
- [4] Claerbout, J., "Migration in Slant-Midpoint Coordinates," SEP 14

FIGURE 1: Synthetic frequency domain source image for examining migration mapping functions. It consists of lines of constant omega. Only half of the k_x values exhibited because of symmetry. All images are 128 by 128 points. $dx = dt = v = 1.0$.



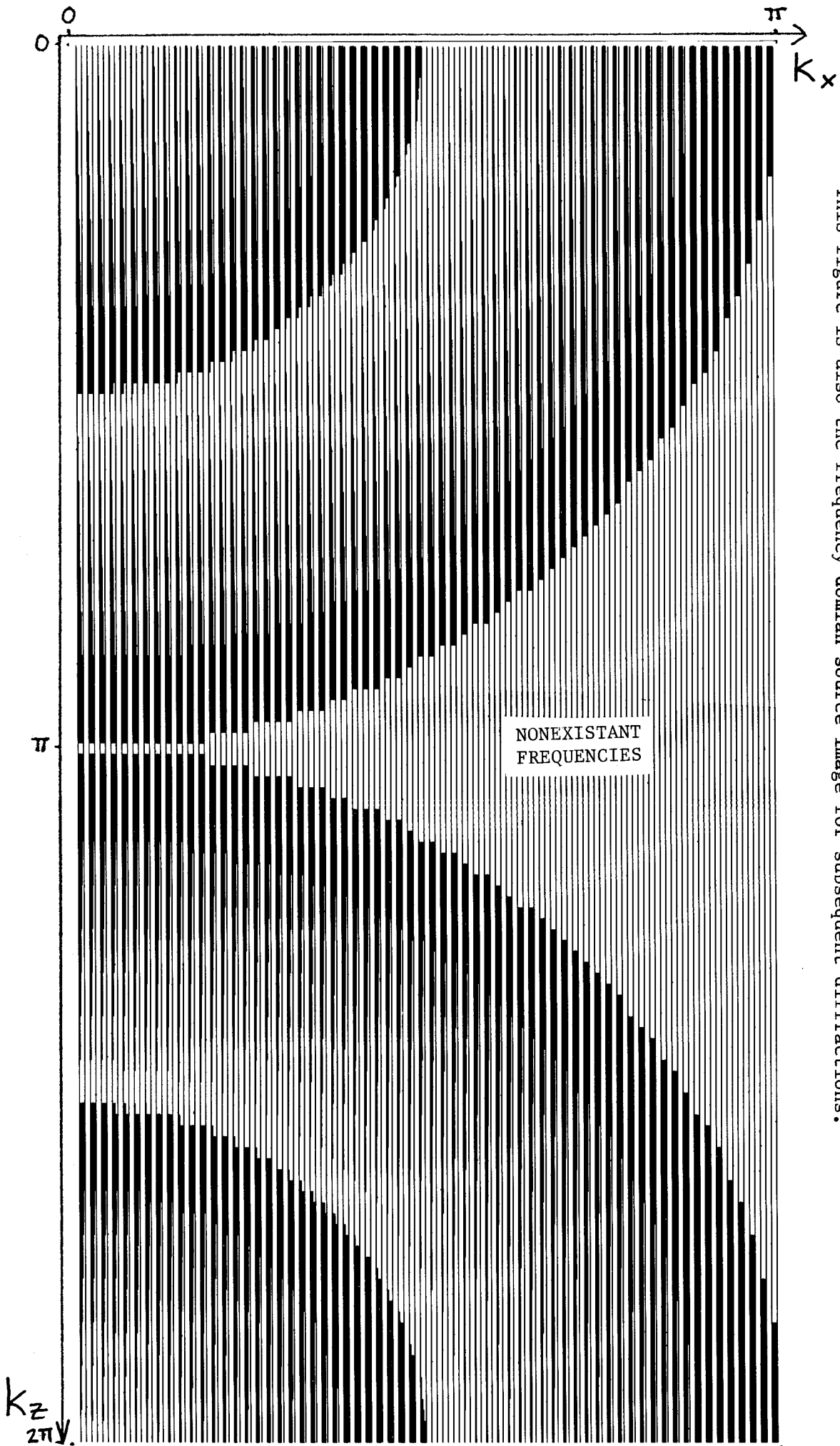


FIGURE 2: Slant-midpoint frequency domain migration mapping of figure 1 for $H = 0$. Lines of constant omega in figure 1 are mapped into circles. This H value is equivalent to CDP stacked migration. Zeroed regions try to map from nonexistent, high omega values of figure 1. This figure is also the frequency domain source image for subsequent diffractions.

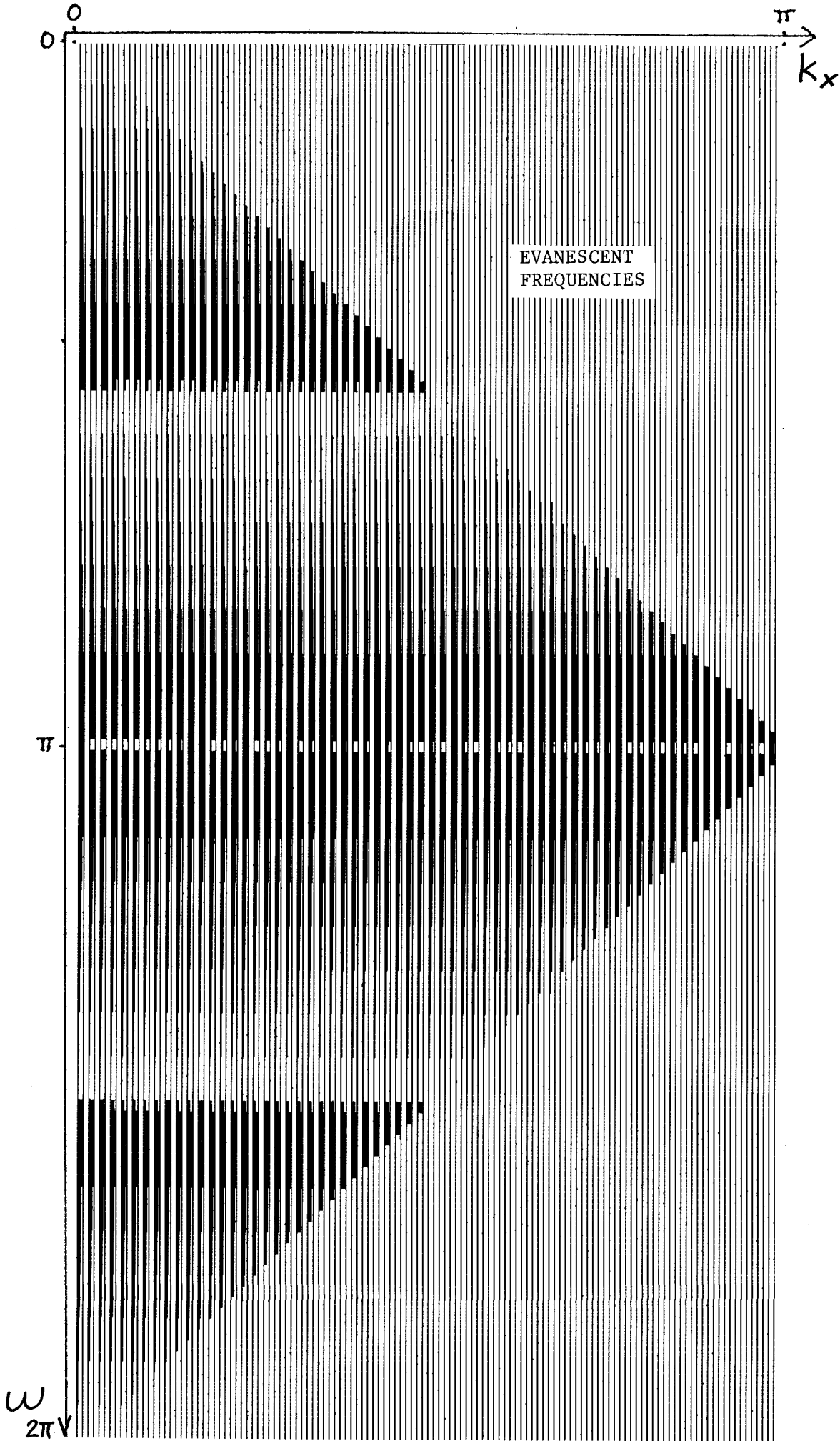


Figure 3: Slant-midpoint frequency domain diffraction mapping of figure 2. Circles have been mapped back into lines. Zeroed regions tried to map from nonexistent, imaginary k_z frequencies of figure 2.

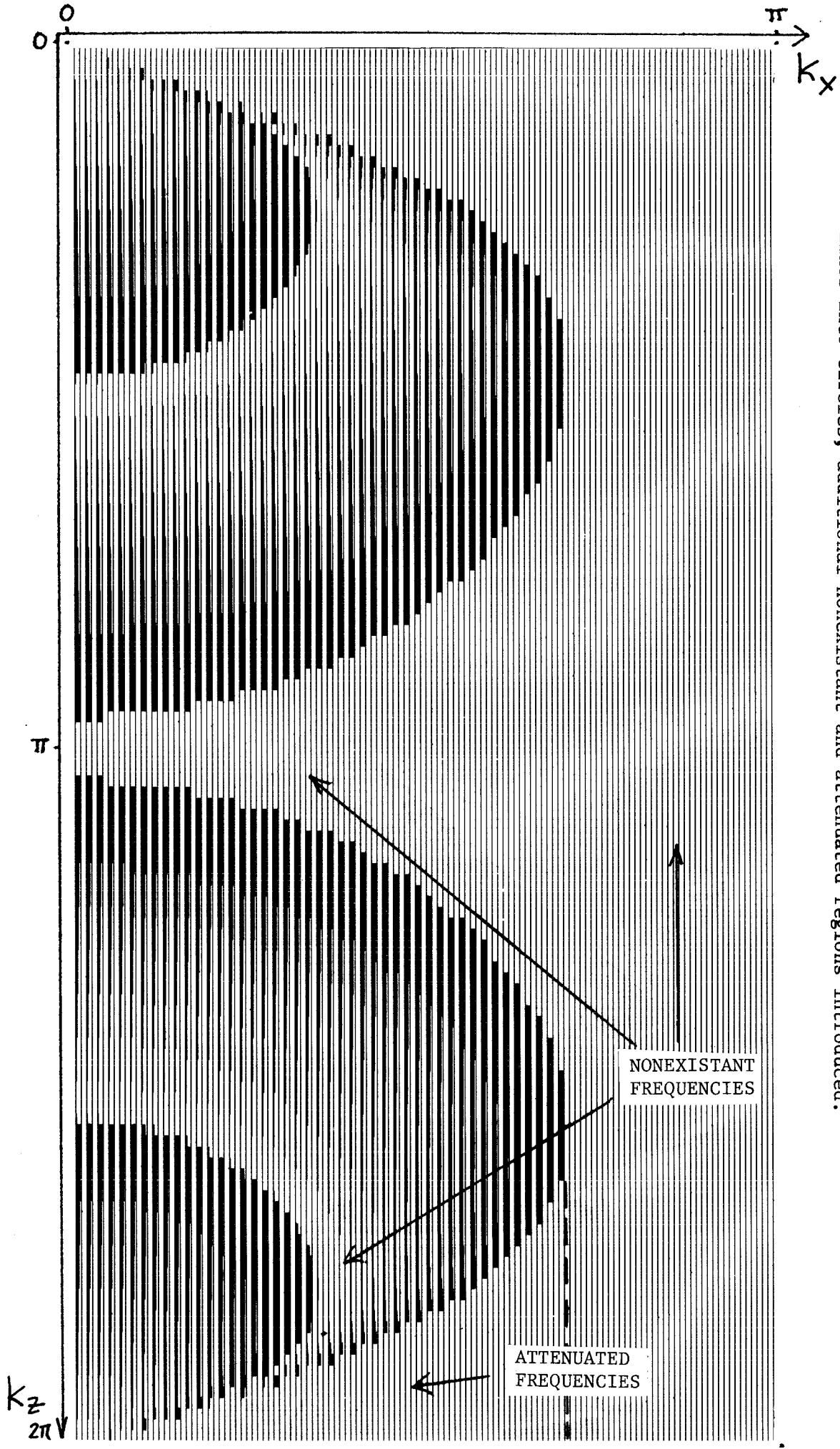


FIGURE 4: Slant-midpoint migration mapping for $H = pv = .3$. No longer the simple mapping of lines into circles, additional nonexistant and attenuated regions introduced.

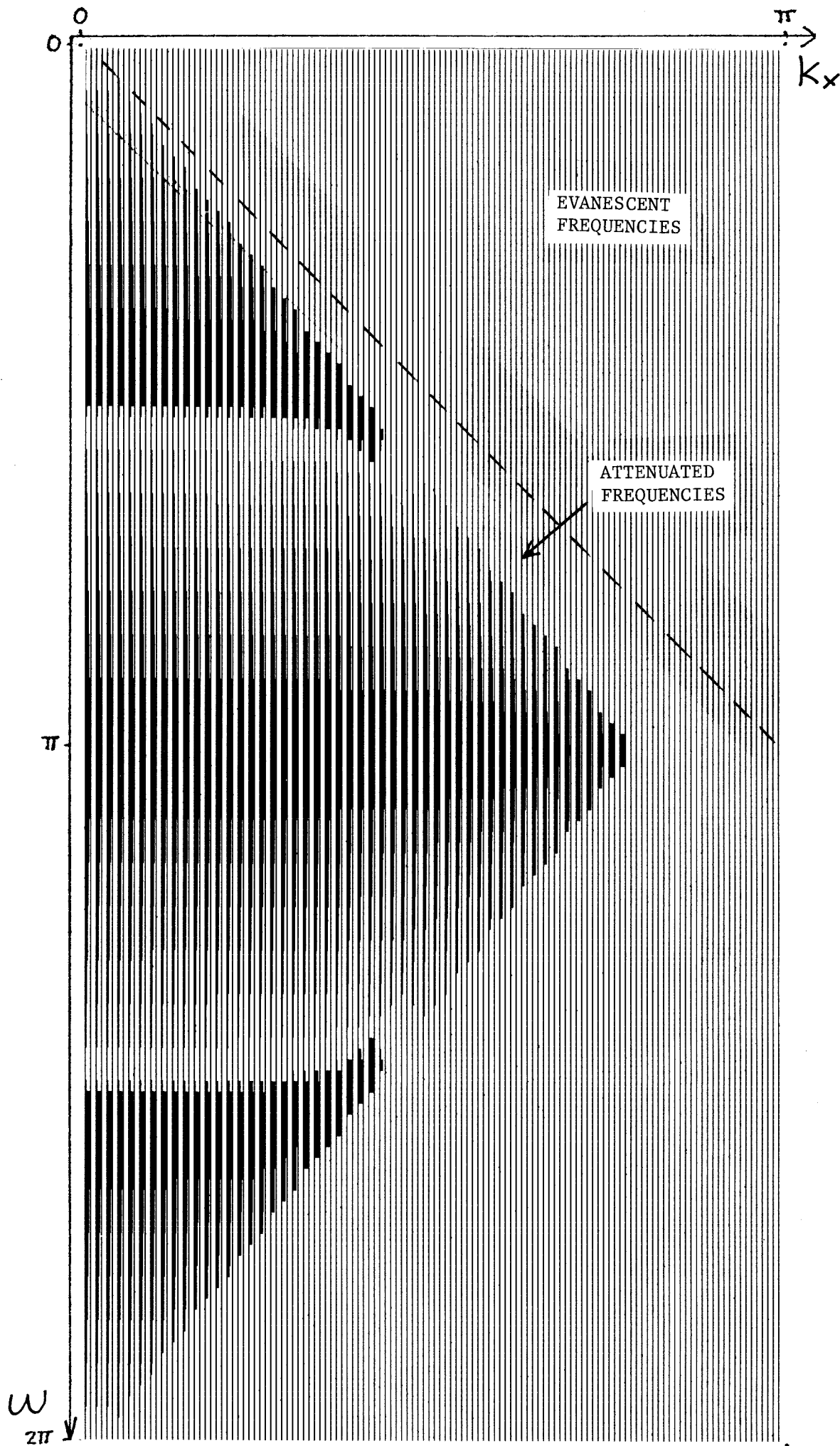


FIGURE 5: Slant-midpoint diffraction mapping for $H = pv = .3$. with attenuated energy zeroed.

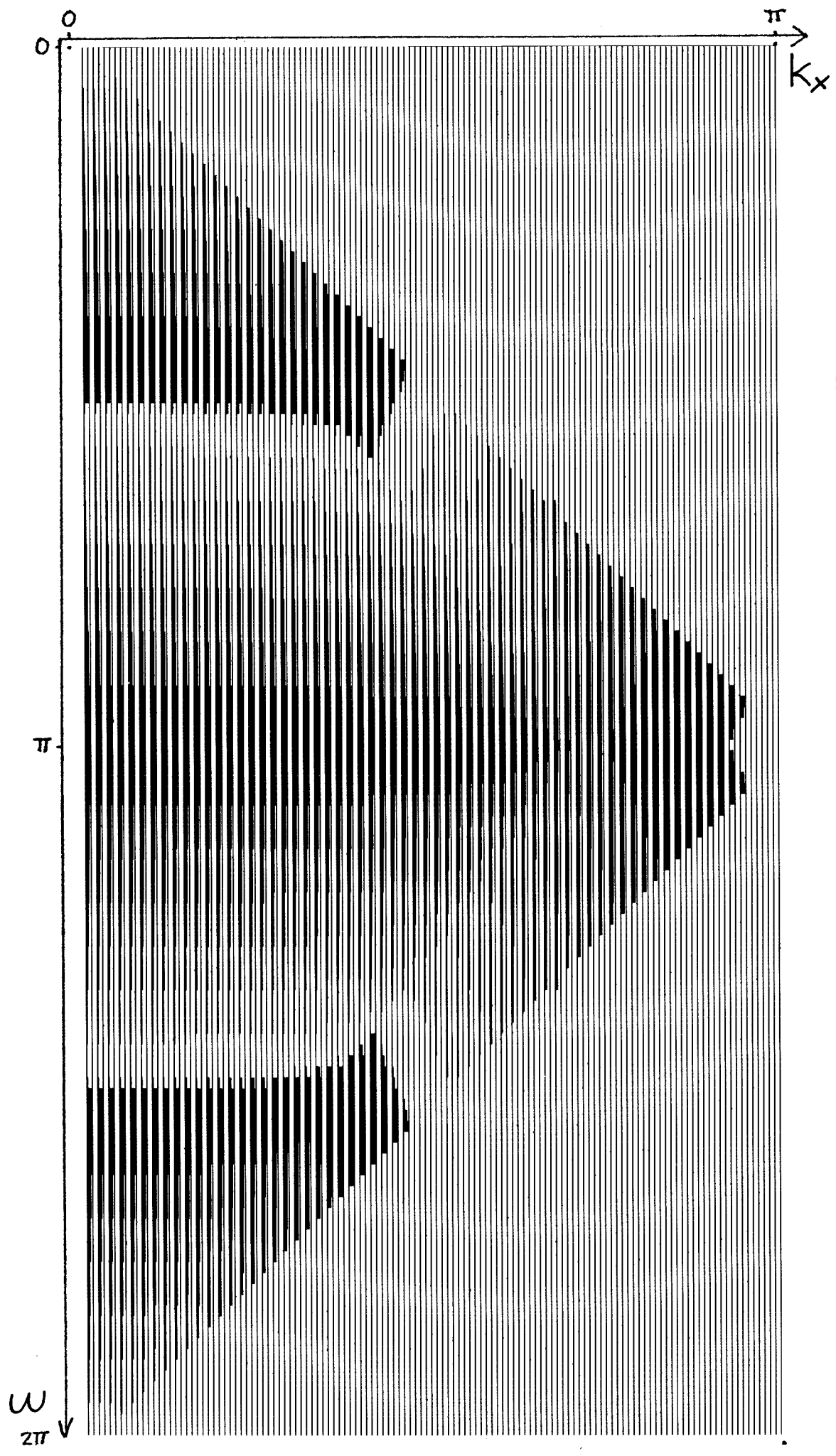


FIGURE 6: Same as figure 5, slant-midpoint diffraction mapping for $H = pv = .3.$, including attenuated energy.

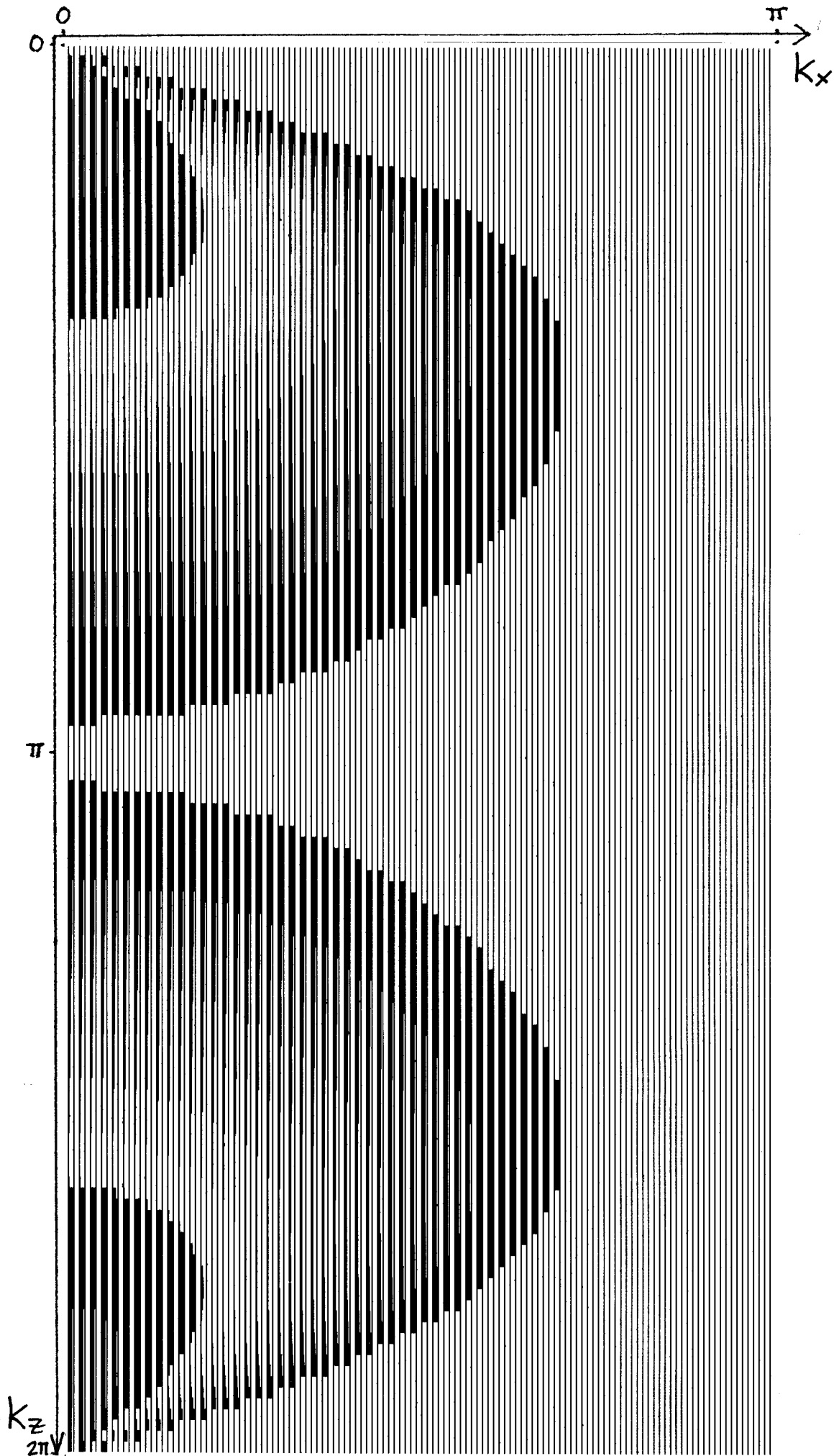


FIGURE 7: Unstacked migration mapping for cross section $kh = .3$. Has superficial resemblance to slant-midpoint migration mapping except considerable amplitude difference at low k_x frequencies.

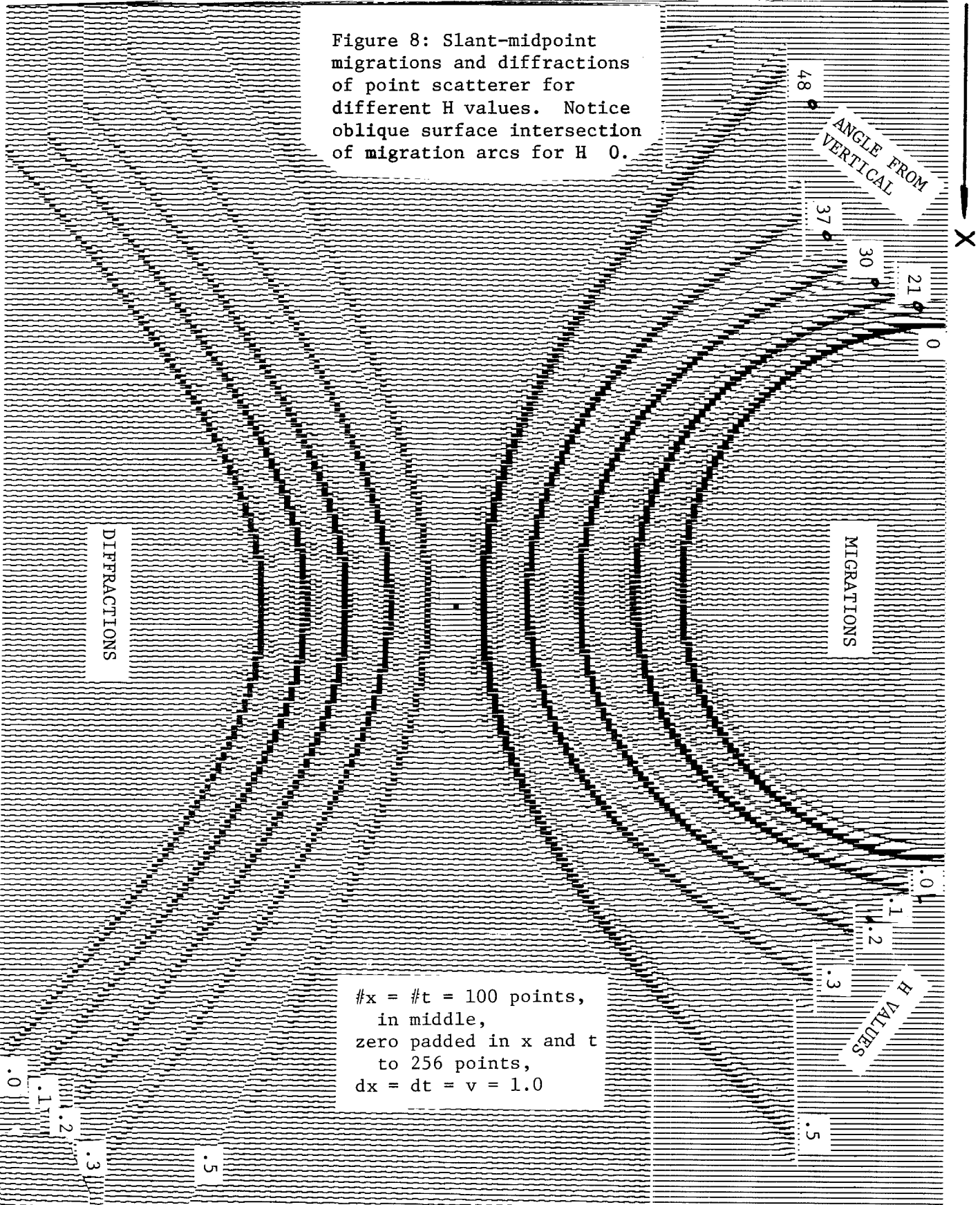
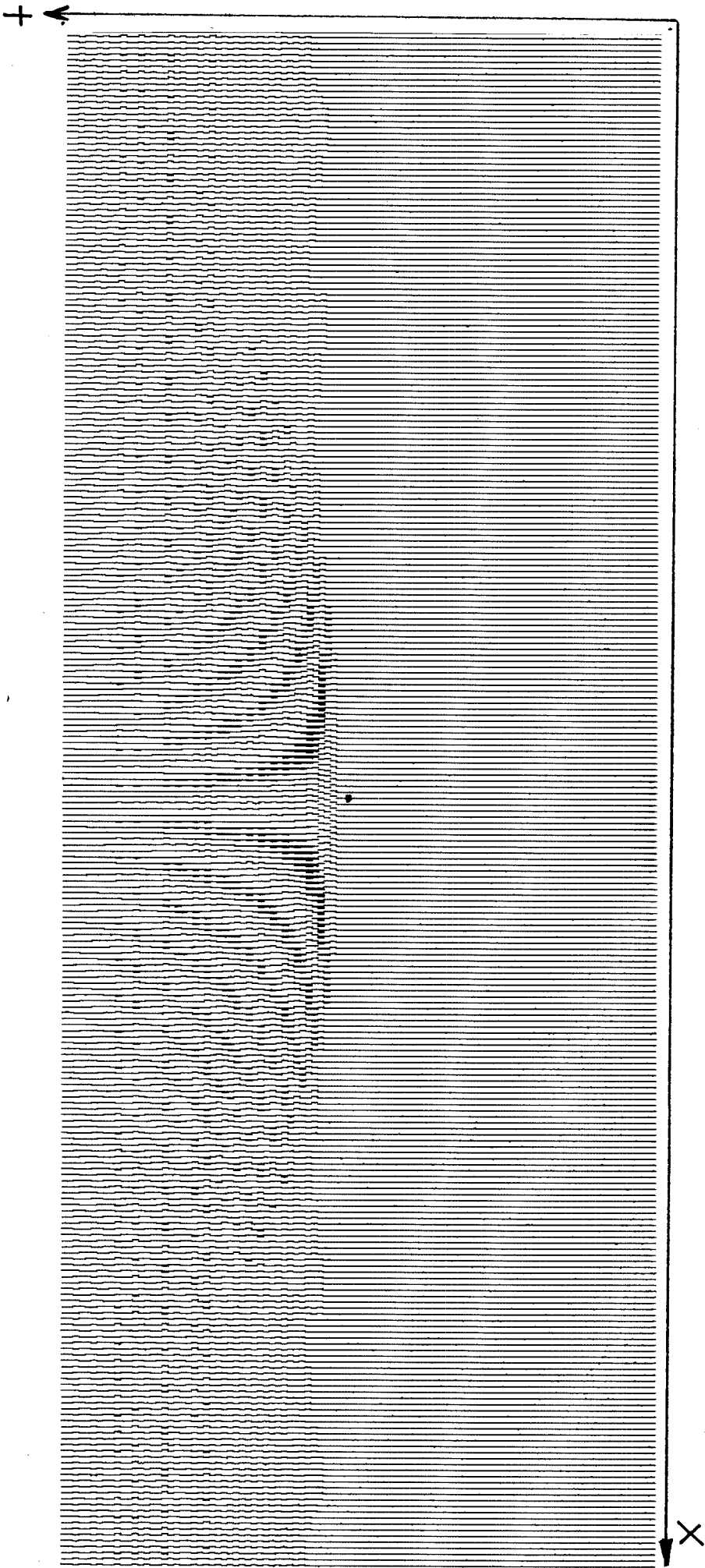


FIGURE 9: Slant-midpoint diffraction of point scatterer for $H = pv = .3$ including attenuated energy.



mapmigh

Author: Ottolini, Stanford Dept of Geophysics, Feb 1978
Machine: PDP 11/70 with FPS array processor, Princeton UNIX fortran

This program is a fast implementation of the frequency domain transfer function for slant-midpoint migration:

$$\text{pmig}(kx, kz) = \frac{\{.5*v*kz**2*(1H**2) - 2*kx**2*kz**2H**2 - kx**4*H**4\}}{\{kx**2 + kz**2\}**1.5 / \{kz**2*(1-H**2) - kx**2H**2\}**1.5} \\ * \text{pdiff}(kx, w = .5*kz* \{kx**2 + kz**2\}**1.5 / \{kz**2*(1-H**2) - kx**2*H**2\}**1.5)$$

Program is middle of 5 step frequency domain migration program set:

- (1) lateral fft $x \Rightarrow kx$
- (2) transpose to time vectors of constant kx
- (3) inner ffts and frequency domain transfer function mapping $t \Rightarrow w$, $w \Rightarrow kz$, $kz \Rightarrow z$
- (4) transpose to kx vectors of constant z
- (5) inverse lateral fft $kx \Rightarrow x$

Program data:

infile: input filename, char*32, transposed lateral fft image,
first $nkx/2+1$ constant kx vectors, remaining vectors
redundant due to real \Rightarrow complex fft symmetry
outfile: output file name, char*32,
 $nkx/2+1$ migrated constant kx vectors
nt: number of time points, integer*2
also number of resulting z points
nkx: number of constant kx vectors, power of 2, integer*2
nw: length of time fft, power of 2, integer*2,
at least 100 greater than nt and less than 4096
dx: depth point separation, real*4
dt: time sample rate in seconds, real*4
v: migration velocity in same units as dx , real*4
h: H value, real*4

Program variables, functions, and subroutines in order of occurrence:

setfil* set device number to file name
nkz number of kz values
nz number of z points
dz z point separation
ixny kx nyquist index
izny kz nyquist index
nw2 w nyquist index / 2
pi 3.14159253
skx, skx2 kx scaling factors
skz kz scaling factor
wny, wny2 w nyquist frequency
h1, h2 constants involving H
ix, xi kx index
cx kx vector
vclr* clear array processor memory
apput* transfer data to array processor
cfft* complex to complex array processor fft
cfftsc* normalize fft result in array processor
apget* transfer data from array processor
kx, kx2 kx wavenumbers
izat, ziat attenuated iz index limit
sqroot modified sqrt zeros negative values


```

      call apget (cx,0,2 * nw,2)
      cx(nw + 1) = cx(1)
c
c      calculate inner loop scaling and index parameters
c      only half of mapping vector need be computed because of symmetry
c      buffering used to overwrite source vector
c
      xi = float (ix - 1)
      kx = xi * skx
      kx2 = kx * kx
      h2 = kx2 * h * h
c      attenuated energy bounds
      ziat = sqrt (h2 / h1) / skz + 1.
      izat = int (ziat)
      if (ziat.ne. float(izat)) izat = izat + 1
      if (izat.gt. izny) izat = izny
c      nonexistant energy bounds
      iznil = int (
c      (sqrt (wny2 * h1 + wny * kx * h - .25 * kx2)
c      + sqrt (wny2 * h1 - wny * kx * h - .25 * kx2)) / skz)+1
      iznil1 = iznil + 1
c
c      first inner loop zeros attenuated kz frequencies
c
      do 5 iz = 1, izat
          cx(iz) = 0.
          cx(nkz - iz + 2) = 0.
5      continue
c
c      second inner loop interpolates existant w locations
c
      if (iznil.lt. izat) goto 50
      do 50 iz = izat, iznil
          zi = float (iz - 1)
          kz2 = zi * zi * skz2
          obliq = v
          cxbuf1 = 0.
          cxbuf2 = 0.
c      mapping index and interpolation weight
          num = kx2 + kz2
          denom = h1 * kz2 - h2
10          if (denom) 40,10,20
          if (zi.ne.0.) goto 40
          zi = 1.
          denom = 1.
20          wi = zi * sqrt (num / denom) + 1.
          iw = int (wi)
          if (iw.gt.nw2) goto 40
          weight = wi - float (iw)
c      obliquity factor
          if (num.eq.0.) goto 40
          obliq = v * ((zi * skz) ** 4. * h1 - (2. * kz2 + h2) * h2)
          / 2. / sqrt (num * denom * denom * denom)
c      map and interpolate
30          cxbuf1 = cintpr (weight,cx(nw-iw+2),cx(nw-iw+1)) * obliq
          cxbuf2 = cintpr (weight,cx(iw),cx(iw + 1)) * obliq
c      overwrite source vector
40          cx(iz) = cxbuf1
          cx(nkz - iz + 2) = cxbuf2
50      continue

```

```

c
c      third inner loop zeros non existant energy
c
c          if (iznil.gt.izny) goto 70
c          do 70 iz = iznil, izny
c              cx(iz) = 0.
c              cx(nkz - iz + 2) = 0.
70      continue
c
c      outer loop inverse ffts (kz=>z) and writes kx strip
c      take vector conjugates before and after fft to correct polarity
c
c          call apput (cx,0,2 * nkz,2)
c          call vneg (1,2,1,2,nkz)
c          call cfft (0,nkz,-1)
c          call vneg (1,2,1,2,nz)
c          call apget (cx,0,2 * nz,2)
c          write (2) (cx(i),i = 1, nz)
80      continue
c          stop
c          end
c
c      functions
c
c      modified square root zeros evanescent values
c
c      function sqroot(arg)
c      if (arg.le.0.) sqroot = 0.
c      if (arg.gt.0.) sqroot = sqrt (arg)
c      return
c      end
c
c      complex geometric interpolation - line 2 (Lynn, SEP 11)
c
c      function cintrp (weight,ca,cb)
c      implicit complex (c)
c      real cabs
c      if (cabs(ca).ne.0.) goto 1
c      cintrp = weight * cb
c      return
1      if (cabs(cb).ne.0.) goto 2
c      cintrp = (1. - weight) * ca
c      return
2      cintrp = ca * cexp (weight * clog (cb / ca))
c      return
c      end

```

```

c      mapdifH
c
c      This program is a fast implementaion of the frequency domain transfer
c      function for slant-midpoint diffraction:
c
c      pdiff(kx,w) = [(1-H**2) + .25*kx*H] / v /
c                  sqrt(w**2*(1-H**2) + w/2/v*kx*H - .25v**2*kx**2)
c                  + [(1-H**2) - .25*kx*H] / v /
c                  sqrt(w**2*(1-H**2) - w/2/v*kx*H - .25v**2*kx**2)]
c      * pmig(kx,kz =
c          [sqrt(w**2/v**2*(1-H**2) + w/2/v*kx*H - .25*kx**2)
c          + sqrt(w**2/v**2*(1-H**2) - w/2/v*kx*H - .25*kx**2)]
c
c      Additional program variables from migration version:
c          iwny          omega nyquist index
c          swv           omega scaling factor / velocity
c          fac1, fac2, fac3  loop constants
c          smap          obliquity factor scaling
c          iwev          evanescent omega index limit
c
c      core of diffraction program:
c
c      calculate outer loop scaling and index parameters
c
c      map into same vector size
c      nt = nz
c      nw = nkz
c      dt = dz / v
c
c      nyquist frequency indices
c      ixny = nkx / 2 + 1
c      iwny = nw / 2 + 1
c
c      compute wave number scale factors from seismogram parameters
c      pi2 = 2. * 3.14159253
c      skx = pi2 / dx / float (nkx)
c      skz = pi2 / dz / float (nkz)
c      swv = .5 * pi2 / v / dt / float (nw)
c      fac1 = swv * swv * (1 - h * h)
c      smap = (1 - h * h) / v
c
c      outer loop reads and ffts (z=>kz) each kx strip
c      only half of the kx strips need be mapped because of symmetry
c
c      do 30 ix = 1, ixny
c          read (1) (cx(i),i = 1,nz)
c          call vclr (0,1,nkz * 2)
c          call apput (cx,0,2 * nz,2)
c          call cfft (0,nkz,1)
c          call cfftsc (0,nkz)
c          call apget (cx,0,2 * nkz,2)
c          cx(nkz + 1) = cx(1)
c
c      calculate inner loop parameters
c
c          kx = float (ix - 1) * skx
c          kx2 = kx * kx
c          fac2 = swv * h * kx
c          fac3 = -.25 * kx2
c          smap1 = smap + .25 * kx * h
c          smap2 = smap - .25 * kx * h
c
c      evanescent index bounds

```

```

c
c first inner loop maps w for nonevanescent kz points
c only half of the mapping vector need be computed because of
c symmetry
c the source vector may be overwritten by the mapped vector by
c going from high to low frequencies and using double buffering
c
      if (iwev.lt.1) goto 15
      cxbuf1 = cx(iwny + 1)
      cxbuf2 = cx(nw - iwny + 1)
      do 10 iw = 1, iwev
          wi = float (iwny - iw)
c mapping index and interpolation weight
          kz1 = sqrt (fac1 * wi * wi + fac2 * wi + fac3)
          kz2 = sqrt (fac1 * wi * wi - fac2 * wi + fac3)
          zi = (kz1 + kz2) / skz + 1.
          iz = int (zi)
          weight = zi - float (iz)
c obliquity factor
          obliq = 1. / v
          if (kz1.gt.0.) obliq = smap1 / kz1
          if (kz2.gt.0.) obliq = obliq + smap2 / kz2
c map and interpolate
          cxbuf3 = cINTRP (weight,cx(nkz-iz+2),cx(nkz-iz+1))*obliq
          cxbuf4 = cINTRP (weight,cx(iz),cx(iz+1))*obliq
c overwrite source vector
          cx(iwny - iw + 2) = cxbuf1
          cx(nw - iwny + iw) = cxbuf2
          cxbuf1 = cxbuf3
          cxbuf2 = cxbuf4
10      continue
          cx(iwny - iwev + 1) = cxbuf3
          cx(nw - iwny + iwev + 1) = cxbuf4
c
c second loop zeros evanescent w points
c
15      iwev = iwny - iwev
          if (iwev.gt.iwny) iwev = iwny
          if (iwev.lt.1) goto 20
          do 20 iw = 1, iwev
              cx(iw) = 0.
              cx(nw - iw + 2) = 0.
20      continue
c
c outer loop inverse ffts (w=>t) and writes kx strip
c take vector conjugates before and after fft to correct polarity
c
          call apput (cx,0,2 * nw,2)
          call vneg (1,2,1,2,nw)
          call cfft (0,nw,-1)
          call vneg (1,2,1,2,nt)
          call apget (cx,0,2 * nt,2)
          write (2) (cx(i),i = 1,nt)
30      continue

```